

# 3 Evolution and modulation of tropical heating from the last glacial 4 maximum through the twenty-first century

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6 Received: 8 November 2010 / Accepted: 28 August 2011  
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8 **Abstract** Twentieth century observations show that during  
9 the last 50 years the sea-surface temperature (SST) of  
10 the tropical oceans has increased by  $\sim 0.5^\circ\text{C}$  and the area of  
11 SST  $>26.5$  and  $28^\circ\text{C}$  (arbitrarily referred to as the oceanic  
12 warm pool: OWP) by 15 and 50% respectively in association  
13 with an increase in green house gas concentrations,  
14 with non-understood natural variability or a combination of  
15 both. Based on CMIP3 projections the OWP is projected to  
16 double during twenty-first century in a moderate  $\text{CO}_2$   
17 forcing scenario (IPCC A1B scenario). However, during the  
18 observational period the area of positive atmospheric  
19 heating (referred to as the dynamic warm pool, DWP), has  
20 remained constant. The threshold SST ( $T_H$ ), which demarks  
21 the region of net heating and cooling, has increased from  
22  $26.6^\circ\text{C}$  in the 1950s to  $27.1^\circ\text{C}$  in the last decade and it is  
23 projected to increase to  $\sim 28.5^\circ\text{C}$  by 2100. Based on climate  
24 model simulations, the area of the DWP is projected to  
25 remain constant during the twenty-first century. Analysis of  
26 the paleoclimate model intercomparison project (PMIP I  
27 and II) simulations for the Last Glacial maximum and the  
28 Mid-Holocene periods show a very similar behaviour, with  
29 a larger OWP in periods of elevated tropical SST, and an  
30 almost constant DWP associated with a varying  $T_H$ . The  
31 constancy of the DWP area, despite shifts in the background  
32 SST, is shown to be the result of a near exact matching  
33 between increases in the integrated convective heating  
34 within the DWP and the integrated radiative cooling outside  
35 the DWP as SST changes. Although the area of the DWP  
36 remains constant, the total tropical atmospheric heating is a  
37 strong function of the SST. For example the net heating has

increased by about 10% from 1950 to 2000 and it is projected  
to increase by a further 20% by 2100. Such changes must be  
compensated by a more vigorous atmospheric circulation, with  
growth in convective heating within the warm pool, and an  
increase of subsiding air and stability outside the convective  
warm pool and an increase of vertical shear at the DWP  
boundaries. This finding is contrary to some conclusions from  
other studies but in accord with others. We discuss the  
similarities and differences at length.

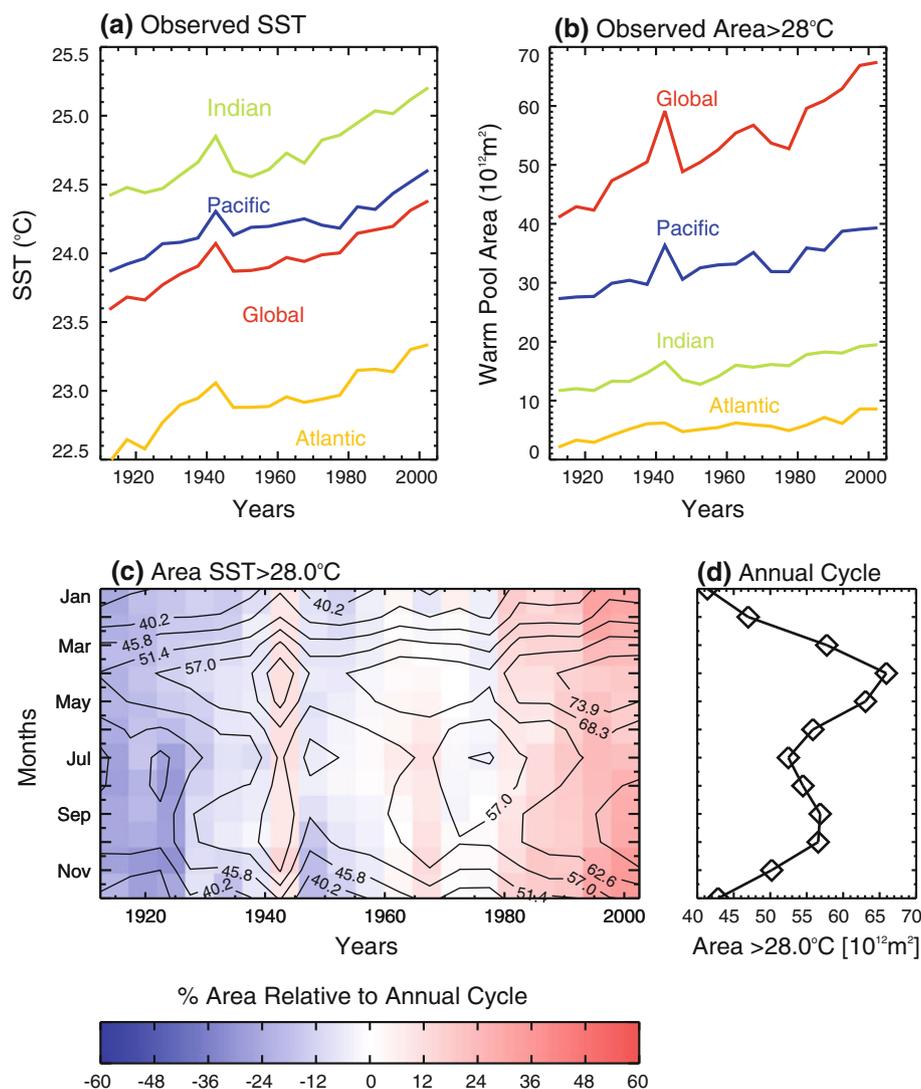
**Keywords** Tropical warm pool · Tropical convective heating · Tropical climate, areas of convection in last glacial maximum, Mid-Holocene and present · Projections of the warm pool for the twenty-first century

## 1 Introduction

A key to understanding the evolution of Earth's climate is  
deciphering the relationship between tropical convection and  
the magnitude and gradients of the local sea-surface temperature  
(SST). Although the variation of SST may not necessarily mirror  
changes in convection, it provides a vivid example of climate  
change. During the twentieth century tropical SST has increased  
by about  $0.8^\circ\text{C}$  in all tropical oceanic basins (Fig. 1a) in two  
major steps. Warming occurred through the 1930s before  
reaching a peak in the 1940–1945 period followed by an SST  
plateau of cooler temperatures for the next 20 years before  
warming rapidly into the present era. Whereas data was scarce  
early in the last century, the trend in global SST has been  
unambiguously positive across the century (e.g., Deser et al.  
2010) whether the trend is caused by increased greenhouse gas  
concentrations, internal climate variability or a combination of  
both.

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**Fig. 1** **a** Evolution of SST for the tropical, Pacific, Indian, and Atlantic basins in 5-year (pentad) bins from 1910 to 2004. **b** Same as **a** but for area ( $10^{12} \text{ m}^2$ ) of SST  $>28^\circ\text{C}$  (OWP area). **c** Evolution of SST  $>28^\circ\text{C}$  ( $10^{12} \text{ m}^2$ ) with the percent deviation relative to the mean annual cycle of area of SST  $>28^\circ\text{C}$ . **d** Long-term mean annual cycle. NOAA Extended Reconstructed SST v2 data was used in this diagram



70 Figure 1b shows the change in area of the ocean warm  
 71 pool (OWP), defined arbitrarily by the  $28^\circ\text{C}$  sea-surface  
 72 temperature (SST) isotherm. We will later abandon this  
 73 arbitrary limit. However, the  $0.8^\circ\text{C}$  increase of tropical SST  
 74 since 1910 has been accompanied by a 70% expansion of  
 75 the area of the OWP (so defined). The behaviour of each  
 76 ocean basin is similar to the overall variation in the global  
 77 tropics (Fig. 1b). The area of the Indian and Pacific OWP  
 78 have increased by about 80 and 50%, respectively, whereas  
 79 the Atlantic OWP, the smallest and coolest of the tropical  
 80 ocean basins, has increased by 350%. To provide a per-  
 81 spective, during an El Niño event the area of the Pacific  
 82 OWP (i.e., the area with SST  $>28^\circ\text{C}$ ) increases by about  
 83 15%. Figure 1c shows a time section of the annual cycle of  
 84 the OWP. The increase in area is spread rather uniformly  
 85 across the annual cycle. The OWP area in the 2000–2004  
 86 pentad is  $\sim 40\%$  larger than in 1980 compared to the cen-  
 87 tury long climatological value. The climatological annual

cycle of the OWP area (Fig. 1d) shows two peaks following  
 the equinoxes with a maximum in April ( $65 \times 10^{12} \text{ m}^2$ )  
 and a secondary peak in October ( $58 \times 10^{12} \text{ m}^2$ ). Based  
 on simulations using climate models forced by a number  
 of emission scenarios (described in Special Report on  
 Emission Scenarios: SRES, IPCC 2000) the increase of  
 the area encompassing  $28^\circ\text{C}$  is projected to continue  
 throughout the twenty-first century (IPCC 2007, e.g.,  
 Knutti et al. 2008).

The general increasing trend in global SSTs during the  
 twentieth century is non-uniform with a peak occurring in  
 the 1930–1940s period and a deep valley between 1950 and  
 1970. There are differing explanations for the very obvious  
 1930–1945 “bump” in SST and OWP area easily identi-  
 fiable in Fig. 1a, and b. Thompson et al. (2008) argue that  
 the bump is the result of an uncorrected instrumental biases  
 in the SST record reflecting a general change from the use  
 of bucket temperature measurements to determine SSTs

106 prior to engine intake measurements generally after 1945.  
 107 However, the land temperature record in the tropics and  
 108 extratropics of both hemispheres and in the Arctic  
 109 (Polyakov et al. 2003) shows a general warming starting in  
 110 the early 1930s up to 1945 and a fall after 1945 through  
 111 1970 with a peak to valley amplitude of about 1°C. The dip  
 112 in surface temperatures is usually attributed to an increase  
 113 in aerosols (IPCC 2007). Irrespective of the cause of the  
 114 mid-century climate event, it does not alter the long-term  
 115 tendencies in SST and OWP area although it may raise  
 116 some questions related to attribution of temperature  
 117 change.

118 Tropical SSTs have also changed on century and mil-  
 119 lennial time scales as suggested by proxy paleoclimate data  
 120 records (e.g., Gagan et al. 2004; Mayewski et al. 2004;  
 121 Furtado et al. 2009). The last glacial maximum (LGM),  
 122 occurring 25–16 ky BP (before present) with the coldest  
 123 period near 21 ky BP, is a key epoch where climate  
 124 remained relatively stationary and considerably cooler than  
 125 present for a prolonged period of time. After 18 ky BP the  
 126 planet started to emerge from a glacial climate of low CO<sub>2</sub>  
 127 concentrations (~200 ppm, Sigman and Boyle 2001) and  
 128 large northern hemisphere ice sheets into the current warm  
 129 Holocene epoch. By the Mid-Holocene period (M-HP, 6 ky  
 130 BP; Mayewski et al. 2004) the continental ice sheets had  
 131 disappeared and CO<sub>2</sub> levels reached pre-industrial period  
 132 (P-IP) levels (~280 ppm, Sigman and Boyle 2001). The  
 133 predominant difference in climate forcing at the time of the  
 134 LGM compared to the present was due to precessional  
 135 changes in Earth's orbit, producing localized changes in the  
 136 distribution of incoming solar radiation (Gagan et al.  
 137 2004). Both the LGM and the M-HP are recognized as  
 138 benchmark periods for climate reconstructions given the  
 139 considerable amount of proxy observations available for  
 140 this period.

141 CLIMAP (Climate Long-Range Investigation, Mapping,  
 142 and Prediction; CLIMAP 1976, CLIMAP 1981) was the  
 143 first international effort focused on the reconstruction of the  
 144 climate state of the surface of the Earth during the LGM.  
 145 According to the CLIMAP reconstructions, tropical SSTs  
 146 were, on average, ~2 K below current values, while the  
 147 extratropics were significantly colder, more so in the  
 148 northern hemisphere than the southern. However, changes  
 149 in tropical warm pool may have been underestimated by  
 150 CLIMAP. Reconciliation of LGM terrestrial data over  
 151 Indonesia with ocean data suggests that the tropical Pacific  
 152 may have been cooler by 2–4°C (Webster and Streten 1978;  
 153 Lea et al. 2000; Rind 2000; Rosenthal et al. 2003; Visser  
 154 et al. 2003; Gagan et al. 2004). In addition, many coupled  
 155 climate models cannot reconcile CLIMAP's relatively  
 156 small changes in tropical sea-surface temperature and have  
 157 simulated larger changes than the CLIMAP reconstructions  
 158 suggest (e.g., Pinot et al. 1999; Otto-Bliesner et al. 2009).

159 Also, the CLIMAP tropical ocean temperatures do not seem  
 160 consistent with the much drier tropical climate suggested by  
 161 New Guinea and North Australian terrestrial data (e.g.,  
 162 Webster and Streten 1978) or tropics-wide aridity (Pinot  
 163 et al. 1999). Furthermore, some studies suggest that during  
 164 the M-HP, SSTs were very similar to present day values,  
 165 except at high northern latitudes and a few coastal regions,  
 166 with differences within the errors associated with recon-  
 167 struction methods (Rahmstorf 2002; Mayewski et al. 2004).  
 168 Other studies suggest that SSTs in the Indo-Pacific oceanic  
 169 warm pool may have been even 1 K higher than present  
 170 values (Gagan et al. 2004). At more distant times in the past,  
 171 Cretaceous tropical SSTs were probably warmer by 4–7 K  
 172 than the present (Norris et al. 2002) but were likely to have  
 173 had the same SST probability density function distribution  
 174 as the current climate albeit shifted to higher temperatures  
 175 (Williams et al. 2009).

176 During the last three decades there has been a continuing  
 177 discussion of whether there is a natural limit to tropical SST.  
 178 Most notable are the thermostat hypotheses first promulgated  
 179 by Newell (1979) and later by Ramanathan and Collins  
 180 (1991). The essence of these theories was the existence of  
 181 negative feedbacks associated with cloud growth and  
 182 increases in evaporation as a result of warmer SSTs. It was  
 183 proposed that these processes kept the warm pool tempera-  
 184 ture within rather narrow bounds. A number of studies  
 185 argued against the existence of such a thermodynamic limit  
 186 (e.g., Wallace 1992; Hartmann and Michelsen 1993;  
 187 Pierrehumbert 1995; Sud et al. 2008; Williams et al. 2009).  
 188 In fact, in the 20 years since the publication of the Newell  
 189 and Ramanathan and Collins papers the temperature of the  
 190 tropical oceans has risen by about 0.3–0.35°C and the area of  
 191 the arbitrary OWP has expanded considerably as can be seen  
 192 clearly in Fig. 1a and b.

193 If organized convection is not bounded by some “trigger  
 194 SST”, so called by Williams et al. (2009), as suggested by  
 195 a number of studies (e.g., Dutton et al. 2000; Sud et al.  
 196 2008) an immediate question arises. What has happened to  
 197 the magnitude and location of convective activity during  
 198 the twentieth century while SST has changed or what was  
 199 its state during previous epochs such as the LGM and the  
 200 M-HP? Furthermore, what may happen to the location and  
 201 intensity of convective activity if the SST were to rise  
 202 throughout the next century?

203 The distribution of tropical convection is probably more  
 204 important from a climate perspective than the absolute  
 205 value of SST itself except perhaps for its influence on  
 206 marine ecosystems systems such as coral reefs and fish  
 207 habitats. The distribution of tropical heating and cooling  
 208 plays a crucial role in global weather and climate vari-  
 209 ability, modulating the spatial distribution of available  
 210 potential energy and providing sources and sinks of inter-  
 211 nal energy. In the tropics, the column integrated heating

(CIH) of an atmospheric column is dominated by latent heat release (e.g., Palmer and Mansfield 1984) while at higher latitudes radiative cooling to space is the dominant factor. In particular, the region of net atmospheric tropical heating, basically the area of organized tropical convection, acts as the “boiler box” of the planetary heat engine (e.g. Palmer and Mansfield 1984; Webster and Lukas 1992; Pierrehumbert 1995). The ascending branches of the Walker and Hadley circulations reside within areas of positive net heating, where precipitation exceeds evaporation and where tropical cyclones form and intensify. The subtropical anticyclonic regions and the descending parts of the Walker and Hadley cells reside within this negative heating region.

The net heat input into the tropical atmosphere, or more specifically, the differential heating between the low and high latitudes, drives the atmospheric circulation that, in turn, promotes wind driven ocean circulations across all spatial and temporal scales advecting net heat poleward and maintaining the global heat balance. The energy transport driven by the dynamical response to global scale differential heating renders the climate at any location to be locally unbalanced. For example, tropical diabatic heating generates large-scale stationary waves and teleconnection patterns (e.g., Bjerknes 1969; Webster 1972; Trenberth et al. 1998) that modulate the meridional heat and momentum transports affecting extratropical weather and climate. This modulation also has consequences for numerical simulations. Erroneous representations of heating in tropical regions force systematic errors in extratropical weather (e.g. Hollingsworth et al. 1980; Simmons 1982).

A number of studies (e.g., Sun and Liu 1996; Seager and Murtugudde 1997; Loschnigg and Webster 2000; Stephens et al. 2004) have argued that the coupled dynamics of the ocean and atmosphere are necessary to determine the location and magnitude of tropical convection. It has also been argued that the SST-rainfall relationship is mostly modulated by large-scale moisture convergence arising from differential heating of the atmosphere between the warm pool and adjacent regions, large-scale subsidence induced by remote convection and the magnitude of the local surface latent heat flux (e.g. Fu et al. 1992; Webster 1994; Emanuel et al. 1994; Pierrehumbert 1995; Lau et al. 1997; Li et al. 2000). This reasoning implies that both the transition between different convective regimes, the onset of deep convection and total tropical CIH are functions not only of the local SST, but also of the three-dimensional large-scale atmospheric circulation.

During the last decade there has been a debate on whether the Hadley and Walker Circulation has decreased in intensity in the present era and will continue the same trend with further global warming (e.g., Vecchi et al. 2006;

Held and Soden 2006; Vecchi and Soden 2007) or has increased and perhaps will continue to increase in intensity in the future (e.g., Mitas and Clement 2005, 2006; Meng et al. 2011; Zhou et al. 2011). These are fundamental questions that are addressed in this paper.

In summary, the SST distribution has varied in the past as well as during the present era, and there is a high probability of changes occurring during the next century. However, a number of fundamental questions still remain unanswered. How has large-scale convection, paramount in the heating of the tropical atmosphere varied in the past and how it will vary in the future? Did the intensity of convection increase, decrease or remain the same as the magnitude of the SST waxed and waned? Did the area of heating expand or contract as the SST patterns changed? Clearly, the location, size and magnitude of the area of positive heating are critical pieces of information that determine the character of the Hadley, Walker and monsoon circulations and teleconnections to higher latitudes. A major thrust of this paper is to sort out the relationships between the SST and regions of positive and negative net heating in order to see if the characteristics of the tropical atmosphere in the past and present can be explained consistently and the future changes anticipated from a physical a priori basis.

In the next section we outline our methodology and describe the data and models used in the study. Results of the study of reanalysis data sets and the output from coupled climate models are presented in Sect. 3. In Sect. 4, simple theoretical and intermediate numerical models are developed to aid in deciphering the results from data analysis and from complex coupled models. A summary of the results is given in Sect. 5.

## 2 Data and methodology

The NOAA extended reconstructed SST data set (Smith and Reynolds 2004) is used to examine twentieth century variations of SST and the OWP in the 1910–2004 period. Although the data set is available from 1854 to the present, tropical SST measurements before 1910, especially over the Pacific Ocean, are scarce (e.g., Deser et al. 2010). The reconstructed SST, initially constructed by Oberhuber (1988), is available globally on a 2° grid and is constructed using the available International Comprehensive Ocean–Atmosphere Data Set SST data at <http://www.ncdc.noaa.gov/oa/climate/coads/>. NCEP–NCAR (Kalnay et al. 1996) and ECMWF ERA-40 (Uppala et al. 2005) reanalysis data sets, both available on a global 2.5° grid and extending from 1948 to present and 1957–2002, respectively, are used to derive fields of CIH. Following Nigam (1994), the three-dimensional heating rate,  $\dot{Q}$ , is calculated as a residual in

315 the thermodynamic equation using the monthly fields of  
316 both reanalyses, i.e.

$$\begin{aligned} \dot{Q} = & \frac{\Delta T}{\Delta t} + \bar{\mathbf{v}} \cdot \nabla \bar{T} + \left(\frac{p}{p_0}\right)^{R/C_p} \bar{\omega} \frac{\partial \bar{\Theta}}{\partial p} \\ & + \left(\frac{p}{p_0}\right)^{R/C_p} \left( \nabla \cdot \overline{\mathbf{v}'\Theta'} + \frac{\partial(\overline{\omega'\Theta'})}{\partial p} \right) \end{aligned} \quad (1)$$

318 where,  $\omega$  is the vertical velocity,  $T$  temperature,  $t$  time,  
319  $p$  pressure,  $R$  the gas constant of air,  $C_p$  the specific heat at  
320 constant pressure and  $\Theta$  the potential temperature. The  
321 overbar denotes a monthly average and prime represents  
322 the departure of the 6-h analysis from the monthly average.  
323 Given the approximate thermodynamic balance in the  
324 tropics between convective heating rate  $\dot{Q}$  and  $\bar{\omega}\partial\bar{\Theta}/\partial p$   
325 (Charney 1969), the heating directly associated with con-  
326 vection  $\dot{Q}_c$  is given by:  $\dot{Q}_c = (p/p_0)^{R/C_p} \bar{\omega}\partial\bar{\Theta}/\partial p$ . Subse-  
327 quently, derivations of the three-dimensional heating  
328 structure will be restricted to the third term in Eq. 1. It is  
329 very difficult to obtain 6-hourly data from the different  
330 model paleoclimate and twentieth and twenty-first century  
331 model simulations whereas calculations of convective  
332 heating are quite straightforward. This simplification will  
333 facilitate intercomparison of reanalysis products and the  
334 output form different models.

335 Three different sets of numerical simulations are also  
336 used to assess changes of the tropical OWP and tropical  
337 heating. These are (1) the NCAR Parallel Climate Model  
338 (PCM) twentieth century simulations, (2) the World Cli-  
339 mate Research Programme's (WCRP) Coupled Model  
340 Intercomparison Project phase 3 (CMIP3) twentieth and  
341 twenty-first century runs and (3) the Paleoclimate Model-  
342 ing Intercomparison Project (PMIP) LGM and M-HP  
343 (M-HP) simulations.

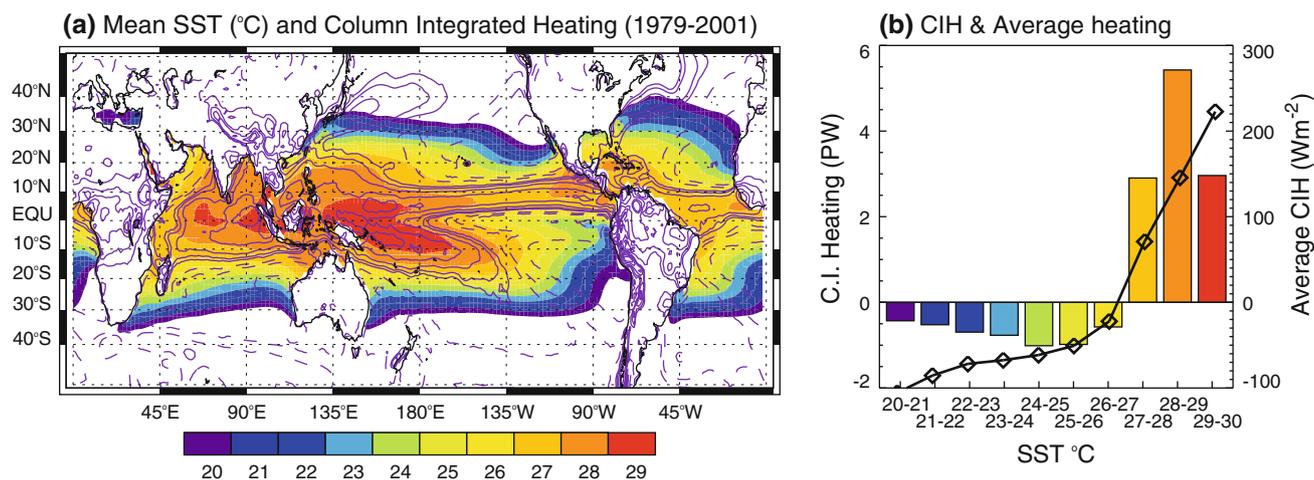
344 Simulations from the World Climate Research Pro-  
345 gramme's (WCRP) Coupled Model Intercomparison Pro-  
346 ject Phase 3 (CMIP3, Meehl et al. 2007) are used to  
347 investigate the variability of twentieth century OWP, the  
348 magnitude and location of the convective threshold tem-  
349 perature  $T_H$ , separating convective from non-convective  
350 regions, and CIH, as well as their potential changes during  
351 the twenty-first century. The analysis include CMIP3  
352 simulations under four different IPCC emissions scenarios  
353 (IPCC 2007): (1) twentieth century simulations with cor-  
354 responding anthropogenic and natural forcing, (2) twenty-  
355 first century simulation with  $\text{CO}_2$  concentrations fixed at  
356 year 2000 ( $\text{CO}_2 \sim 360$  ppm, COMMIT scenario), (3)  
357 twenty-first century simulation with  $\text{CO}_2$  concentrations  
358  $\sim 700$  ppm by 2100 (A1B scenario), and (4) twenty-first  
359 century simulation with  $\text{CO}_2$  concentrations  $\sim 820$  ppm by  
360 2100 (A2 scenario). The CMIP3 dataset was obtained from  
361 the Program for Climate Model Diagnosis and

Intercomparison (PCMDI; <http://www-pcmdi.llnl.gov/>)  
through the WCRP CMIP3 site (<https://esg.llnl.gov:8443/>).

362  
363  
364 The coupled simulations of PMIP I and II (Jousaume  
365 et al. 1999; Pinot et al. 1999; Pausata et al. 2009) are used  
366 to explore the state of the OWP and tropical heating during  
367 the LGM and M-HP. Although the solar forcing is not at a  
368 maximum during the M-HP (Mayewski et al. 2004), this  
369 period is preferred in PMIP simulations over the earlier  
370 Holocene since it is not influenced by remnant continental  
371 ice sheets from the previous glaciation. PMIP simulations  
372 consist of atmospheric-only and coupled simulations using  
373 a set of boundary conditions prepared to represent LGM  
374 and M-HP climates. In particular, external forcing includes  
375 the corresponding  $\text{CO}_2$  concentrations for each period,  
376 orbital parameters resulting in insolation changes and ice  
377 sheet extension. For the LGM experiment, ice sheet extent  
378 and height is provided by Peltier (1994) and the atmo-  
379 spheric  $\text{CO}_2$  concentration is prescribed to be 200 ppm as  
380 inferred from Antarctic ice cores (Raynaud et al. 1993;  
381 Sigman and Boyle 2001). Earth's orbital parameters were  
382 changed according to their values at 21 ky BP, but the  
383 insolation changes were small for this time period. For the  
384 M-HP the  $\text{CO}_2$  concentration was prescribed at the P-IP  
385 value of 280 ppm (Raynaud et al. 1993; Sigman and Boyle  
386 2001). The main change in the incoming solar radiation in  
387 this period is due to the displacement of the longitude of  
388 the perihelion (Berger 1978), which intensifies the seasonal  
389 north-south gradient of incoming solar radiation. SSTs are  
390 either prescribed using CLIMAP for LGM and present day  
391 SST for the M-HP, or computed by coupling the atmo-  
392 spheric model to a slab mixed-layer ocean.

### 3 Results

393  
394 The distribution and magnitude of tropical latent heat  
395 release in deep convection is, in any given period, strongly  
396 associated with the distribution of SST. Figure 2a shows  
397 the long-term average (1979–2001) SST distribution with  
398 contours of column-integrated atmospheric heating (CIH)  
399 for the same period superimposed. Warm waters straddle  
400 the equator, extending from the Indian Ocean, across the  
401 Indonesian Archipelago into the western Pacific with a  
402 secondary area in the eastern Pacific crossing Central  
403 America into the Caribbean and the central Atlantic Ocean.  
404 The CIH patterns have a distinct SST signature: contours of  
405 CIH closely follow those of SST with regions of net  
406 positive CIH confined to areas of warmest ocean water.  
407 Figure 2b shows the total tropical CIH (PW:  $10^{15}$  W)  
408 between  $30^\circ\text{S}$  and  $30^\circ\text{N}$ . Spatially averaged CIH ( $\text{Wm}^{-2}$ ),  
409 plotted in  $1^\circ\text{C}$  SST bins for the period 1979–2001. For this  
410 period, heating of the atmospheric column occurs within  
411 the  $27^\circ\text{C}$  SST isotherm, which is the threshold SST



**Fig. 2** **a** Long-term annual mean SST distribution (1979–2001) using NOAA Extended Reconstructed SST v2 data. The continuous *black* contours represent positive column integrated heating (CIH), the *dashed* contours negative CIH. Contours are plotted for  $\pm 20, 50, 100, 200,$  and  $300 \text{ Wm}^{-2}$ . The column integrated heating is from the ERA-40 Atlas computed using the modelled net temperature tendencies.

**b** Total tropical atmosphere CIH in the  $30^{\circ}\text{S}$ – $30^{\circ}\text{N}$  band plotted in  $1^{\circ}\text{C}$  SST bins. The *black line* shows the average column integrated heating in each SST interval. 25, 48 and 27% of the total heating occurs between the 27–28, 28–29 and 29–30°C SST contours, respectively

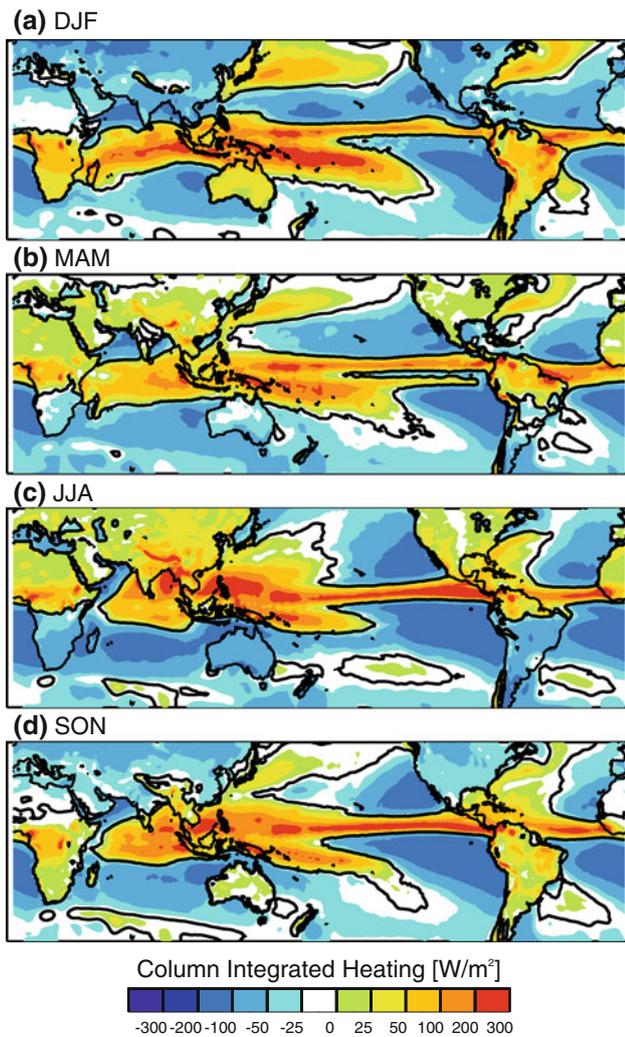
412 bounding the area of large-scale organized tropical convection from adjacent regions of cooling. More intense  
 413 heating, represented by values of spatially averaged CIH  
 414 greater than  $150 \text{ Wm}^{-2}$ , occurs in regions, where SST  
 415  $> 28^{\circ}\text{C}$  (Fig. 2b). Traditionally, the  $28^{\circ}\text{C}$  temperature  
 416 threshold has been used to define the oceanic warm pool  
 417 (OWP), which in the present era covers nearly 30% of the  
 418 global ocean surface. Summing the total vertically integrated  
 419 heating over the ocean in  $1^{\circ}\text{C}$  SST bins shows that,  
 420 on average, 25, 48 and 27% of the total heating occurs  
 421 currently in the 27–28, 28–29 and 29–30°C SST bins,  
 422 respectively. Outside of the region where  $\text{SST} < T_H$ , there  
 423 are broad areas of negative CIH resulting principally from  
 424 long-wave radiative cooling.  
 425

426 Figure 3 shows the seasonal variation of the patterns of  
 427 CIH for the period 1957–2001. A global analysis of CIH  
 428 was performed by Chan and Nigam (2009) although for a  
 429 somewhat shorter period. Their main emphasis was to  
 430 compare the CIH calculated from different reanalysis data  
 431 sets with the heating profiles obtained by Tao et al. (2001)  
 432 based on data from the Tropical Rainfall Measurement  
 433 Mission. Our emphasis is different in that we seek to  
 434 identify relationships of the CIH with the evolving SST  
 435 distribution. Although, most atmospheric heating occurs  
 436 over the OWP there are small regions of heating off the  
 437 extratropical coasts in the northwest Pacific and Atlantic  
 438 due primarily to the heating of cold continental air flowing  
 439 over the warm ocean western boundary currents during  
 440 winter, and from elevated heating from large mountain  
 441 ranges during summer. Specifics of the regional details will  
 442 be the subject of another study. Here we are more inter-  
 443 ested in the broad scale features of the CIH distribution.

Figure 3 also shows that the spatial structure of the atmo-  
 spheric heating varies considerably from season to season  
 in close association with the evolution of SST and solar  
 heating. During the boreal winter, heating is concentrated  
 over the OWP and the southern hemisphere continental  
 areas in addition to heating over the northern hemisphere  
 western boundary currents. During the boreal summer  
 atmospheric heating occurs over the OWP differing spa-  
 tially to winter patterns, particularly in the Indian Ocean  
 and the South Pacific Convergence Zone, and over the  
 northern hemisphere land masses. The boreal summer  
 heating over the western boundary currents is reduced in  
 magnitude and extent compared to that during boreal  
 winter.

### 3.1 Twentieth and twenty-first century OWP

Analysis of twentieth century ocean–atmosphere coupled  
 simulations for different combinations of anthropogenic  
 and natural forcing (Meehl et al. 2004) indicates that the  
 most likely cause of the OWP expansion is the enhanced  
 concentration of atmospheric  $\text{CO}_2$ . Increased  $\text{CO}_2$  con-  
 centration during the twentieth century explains about 60%  
 of the observed increase in the area of the OWP, far larger  
 than the impacts of volcanoes, aerosols and etc. Specifi-  
 cally, the PCM simulations only show an expanding OWP  
 area when green house gas forcing is included in the total  
 forcing. At the same time, it should be noted that internal  
 variability in many models is smaller than observed and  
 that there is less skill in simulating oscillations of the  
 couple ocean–atmosphere system than the variability  
 associated with external forcing.



**Fig. 3** Annual cycle of CIH: **a** December, January and February; **b** March, April and May; **c** June, July and August; and **d** September, October, and November from ERA-40 Atlas for the 1957–2001 period

474 Figure 4 shows the evolution, in 5-year blocks, of the  
 475 uncorrected tropical OWP area from 1920 to 2000, as  
 476 simulated by eight different ocean–atmosphere coupled  
 477 models from the CMIP3 dataset.<sup>1</sup> The relationship between  
 478 of the observed and simulated OWP area relative to their  
 479 corresponding long-term means are also shown as scatter  
 480 plots, including their cross-correlations. Whilst the mag-  
 481 nitude of the OWP area is different between models, owing  
 482 to their different resolutions, all models simulate the  
 483 observed OWP expansion during the twentieth century  
 484 remarkably well with cross-correlations ranging from 0.68

1FL01  
 1FL02  
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 1FL04

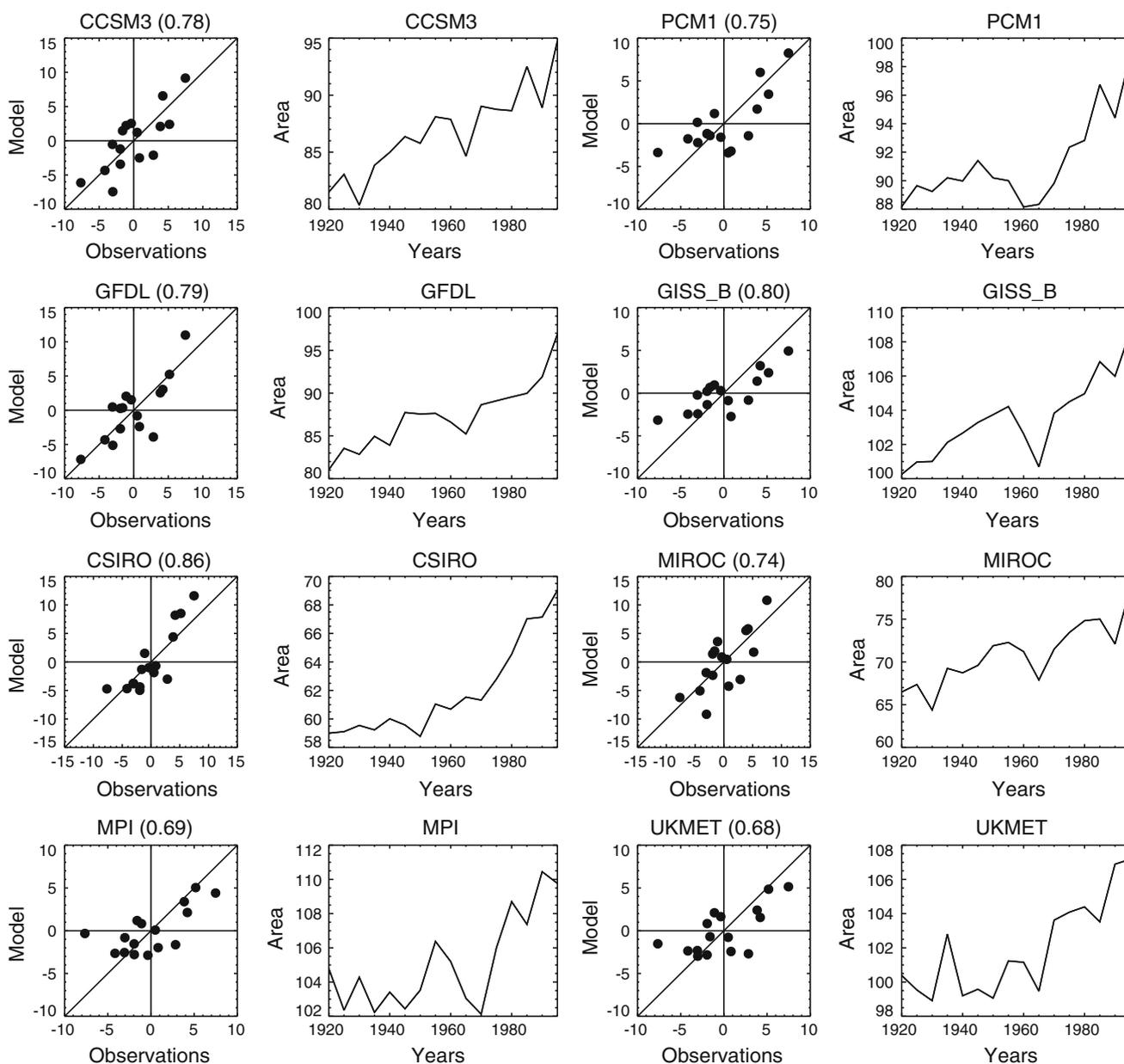
<sup>1</sup> The models that were selected are NCAR CCSM3, NCAR PCM1, NASA GFDL-CM2.1, NASA GISS, MIROC3.2-MEDRES, ECHAM5/MPI-OM, UKMET-HadCM3, CSIRO-Mk3.5. Details of the models and the CMIP3 process can be found at: <https://esg.llnl.gov:8443/>.

to 0.86. Correlations between the observed and the simu-  
 lated area of SST >26.5°C are also high (~0.8 correla-  
 tions). All correlations are significant at a level greater than  
 99%.

In order to account for model biases and to obtain a  
 multi-model estimation of the OWP (area >28°C) and the  
 area where SST >26.5°C during the twentieth century, as  
 well as the projected changes during the twenty-first cen-  
 tury, the simulated OWP and area with SST >26.5°C are  
 statistically corrected relative to the twentieth century  
 observations. Specifically, linear least square regressions  
 between the observed and the simulated OWP area for each  
 model during the twentieth century are used to adjust the  
 magnitude of model biases statistically during the twentieth  
 and twenty-first century. Note though, all models reproduce  
 the observed increase in OWP reasonably well before the  
 bias correction (Fig. 4). All corrections, introduced by the  
 linear fitting equations, only account for the magnitude of  
 the increase and not for the sign of the tendencies. Figure 5  
 shows the bias-corrected changes in OWP and area of SST  
 >26.5°C from the CMIP3 simulations for different emis-  
 sion scenarios. The figure includes the multi-model bias-  
 corrected ensemble mean and the variability estimated as  
 ±1 standard deviation of the different bias-corrected model  
 simulations (see inset, Fig. 5). The correlation between the  
 observed and multi-model simulated OWP area is 0.82 for  
 the 1920–2000 period. This correlation is significant at a  
 level greater than 99%. Projecting into the future, and  
 considering the COMMIT scenario, the OWP and the area  
 of SST >26.5°C stabilizes in about 30–40 years to a level  
 around 10–20% greater than in 2000–2004. Under scenarios  
 A1B and A2, the OWP area continues to increase during the  
 twenty-first century to a multi-model average of +70 and  
 +90% by 2100, respectively, relative to the 2000–2004  
 value. For the area of SST >26.5°C, the increase is +30  
 and +40%, respectively, for scenarios A1B and A2.

The spatial distribution of the observed OWP area for  
 the 1920–1925, 1960–1965, 2000–2005 pentads is shown  
 in Fig. 6. In addition, the projected areas for 2030–2035,  
 2060–2065 and 2090–2095 pentads are also shown for the  
 A1B scenario (Fig. 6). The contours for the twenty-first  
 century OWP area are determined by projecting the relative  
 area changes from the multi-model ensemble onto the  
 2000–2005 average SST with the assumption that the  
 expansion follows the observed spatial SST gradients  
 homogeneously as it did quite closely during the twentieth  
 century. The largest relative changes during the twenty-first  
 century occur in the Atlantic Ocean and the eastern Pacific,  
 while smaller changes are seen in the western Pacific. By  
 2060, the OWP in the Atlantic spans from Central and  
 South America to Africa throughout the entire annual  
 cycle, the western Pacific area merges with the Eastern  
 Pacific during all seasons, and in the eastern Pacific this

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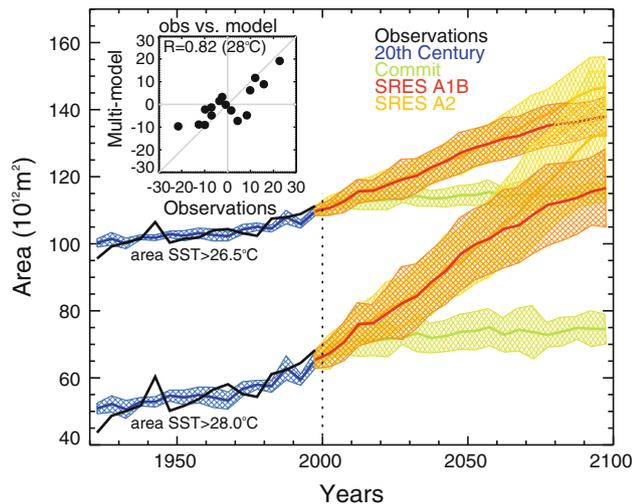
**Fig. 4** Observed twentieth century changes in area of SST >28°C between 1920 and 2000 versus the changes in area in each of the selected CMIP3 models for the twentieth century between 1920 and 2000. Models used are listed in the text

538 region is no longer confined to the northern hemisphere  
 539 during spring. Changes in the Indian Ocean are particularly  
 540 large for fall and winter seasons: by year 2060, the fall  
 541 OWP spans across the entire Indian Ocean from Africa to  
 542 Indonesia.

543 3.2 Twentieth and twenty-first century CIH and its  
 544 relationship with SST: the dynamic warm pool

545 The 50-year evolution of the CIH as a function of SST  
 546 between 20°S and 20°N is computed using the NCEP-  
 547 NCAR and ECMWF ERA-40 reanalyses. Figure 7a shows

the co-evolution of the SST and total CIH from 1950 to 2004. The relationship is non-stationary. The CIH profiles  
 548 during spring. Changes in the Indian Ocean are particularly  
 549 large for fall and winter seasons: by year 2060, the fall  
 550 OWP spans across the entire Indian Ocean from Africa to  
 551 Indonesia. During the last 50 years,  $T_H$  increased almost  
 552 linearly with the average tropical SST about 0.5°C from  
 553 26.6 to 27.1°C (Fig. 7b). In addition, there has  
 554 been ~ 15% increase in net heating within the areas of  
 555 convection during the last 50 years (Fig. 7b). Similar  
 556 results are obtained using ECMWF ERA-40 reanalysis  
 557 (Fig. 7c, d). There are some differences between the NCEP  
 558 and ERA-40 estimations although the progression of  $T_H$   
 559



**Fig. 5** Changes in area of SST  $>28$  and  $26.5^{\circ}\text{C}$  from observations (black) and from selected CMIP3 coupled models for the twentieth century runs (blue) and for the following scenarios: COMMIT (green), A1B (red), and A2 (orange) scenarios. The inset shows the relationship between the observed twentieth century changes in area of SST  $>28^{\circ}\text{C}$  versus the relative area changes in the multi-model ensemble mean of the selected WCRP CMIP3 models for the twentieth century for the 16 pentads between 1920 and 2000. Quantities shown are deviations from long-term mean

560 and the trends in net heating are very similar. Chan and  
561 Nigam (2009) also found differences between the two  
562 estimates of CIH especially across Indonesia. Our biggest  
563 difference occurs around 1970, a period that was not  
564 contained within the Chan-Nigam analysis.

565 Estimates of future changes in CIH under different  
566 emission scenarios are evaluated using the CMIP3 simu-  
567 lations. Figure 8 shows time-SST sections of tropical CIH  
568 between  $20^{\circ}\text{S}$  and  $20^{\circ}\text{N}$  for the twentieth and the twenty-  
569 first centuries (A1B scenario) for four different CMIP3  
570 models. Results from other CMIP3 models are similar, with  
571 all models indicating that  $T_H$  increases with time as SST  
572 increases. Overall, by 2100,  $T_H$  would be about  $1.6^{\circ}\text{C}$   
573 degrees higher than the present day value of  $\sim 27^{\circ}\text{C}$ .

574 Our findings are in agreement with the results of other  
575 studies (e.g. Lau et al. 1997; Dutton et al. 2000; Williams  
576 et al. 2008; Sud et al. 2008). Simply, given the steadily  
577 increasing threshold  $T_H$ , the existence of a fixed convective  
578 threshold is not evident in an evolving climate nor does  
579 such a constant  $T_H$  have, as we will develop, a reasonable  
580 physical basis.

581 Figure 9a shows the area of positive CIH (heavy black  
582 line: termed DWP or dynamic warm pool for later refer-  
583 ence) during the twentieth century and for the IPCC A1B  
584 emission scenario in the twenty-first century for four dif-  
585 ferent CMIP3 models. In spite of the expansion of the areas  
586 above different static SST thresholds, the area of positive  
587 heating region does not show any significant long-term

trends, covering about 25% of the global ocean or about  
588 60% of the ocean area between  $20^{\circ}\text{S}$  and  $20^{\circ}\text{N}$ . The area of  
589 the OWP for each of the scenarios is shown for compari-  
590 son. Similar results are obtained for COMMIT and IPCC  
591 A2 scenarios (not shown). Thus, while the SST and the  
592 area of the OWP area increase sharply with time, the area  
593 of net heating in the tropics (the DWP), remains remark-  
594 ably constant over the twentieth and twenty-first century.  
595

Dutton et al. (2000) also noted a constancy of the area of  
596 deep tropical convection defined by out-going longwave  
597 radiation and increasing threshold temperatures as  $\text{CO}_2$   
598 concentrations increase in a set of coupled general circula-  
599 tion model experiments. Rather than a static definition set  
600 by a constant temperature, the climatically active warm  
601 pool is defined dynamically by the large-scale coupled  
602 ocean-atmosphere system, and as previously stated, cor-  
603 responds to the area where SST  $> T_H$  enclosing the area of  
604 positive CIH.  
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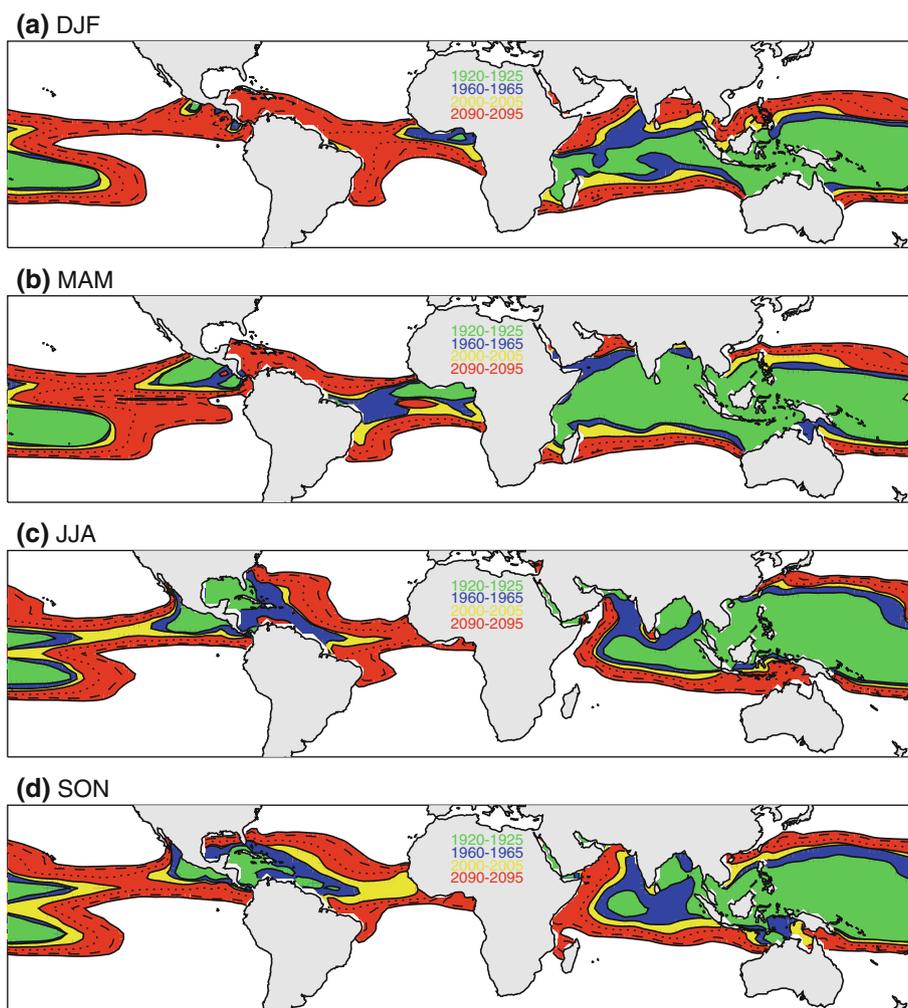
Figure 9b shows the change in CIH relative to the year  
606 2000 value in each of the models. Despite the near con-  
607 stancy of the area of positive CIH, the magnitude of  
608 heating within the area increases substantially in the  
609 twentieth and twenty-first century simulations. The  
610 increase during the twentieth century is similar to that  
611 calculated from the reanalysis products (Fig. 7). The  
612 change in magnitude of the net heating is particularly sharp  
613 during the twenty-first century with a 20% increase relative  
614 to the simulated 1995–2000 values. Such projected  
615 increase in CIH, together with a constant area of positive  
616 heating would suggest a more vigorous circulation, with  
617 intensified moist convection in tropical areas and enhanced  
618 subsidence outside the tropics, also leading to a stronger  
619 vertical wind shear nearby the convective region's  
620 boundaries. The intensified hydrological cycle could have  
621 potentially important thermodynamic and dynamic impacts  
622 for the tropical as well as the global climate.  
623

In Sect. 4, we discuss how a combination of thermo-  
624 dynamic and dynamic factors the location and size of the  
625 DWP to fall within a relatively narrow range.  
626

### 3.3 Paleo-CIH/SST distributions and relationships 627

628 Significant changes in climate are not exclusive to the  
629 present era and are not necessarily induced by human  
630 activities. As mentioned earlier, the LGM and the M-HP are  
631 two periods where the Earth's climate was approximately in  
632 balance for long periods and significantly different from  
633 today's climate. These two periods constitute an opportunity  
634 to evaluate whether or not twentieth and twenty-first century  
635 changes in OWP and positive CIH area (or DWP) are con-  
636 sistent with past climates under very different conditions. We  
637 use the PMIP I and II simulations of the LGM, M-HP and  
638 P-IP to estimate changes in the DWP and the magnitude of

**Fig. 6** Spatial distribution of the area of SST  $>28^{\circ}\text{C}$  for the 1920–1925 (green), 1960–1965 (blue), 2000–2005 (yellow) pentads from NOAA Extended Reconstructed SST v2 data as well as the 2030–2035 (dotted line), 2060–2065 (dashed line) and 2090–2095 (red) pentads from the A1B scenario



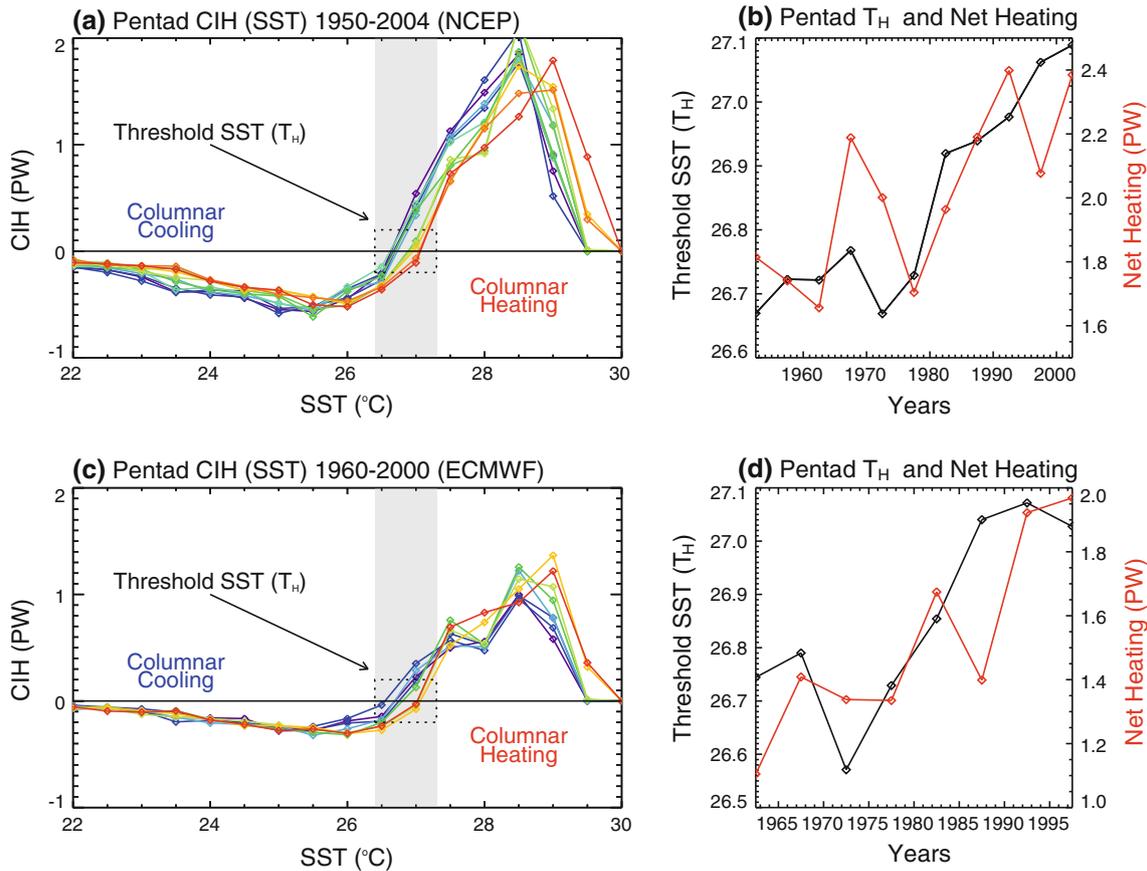
639 convection and heating within the DWP, relative to changes  
640 in tropical SST.

641 Figure 10 presents a summary of the analysis of four  
642 different PMIP I models (NCAR CCM3, Canadian Climate  
643 Center model, GFDL, and UGAMP) for the LGM and the  
644 P-IP. The PMIP I data that are publically available are not  
645 sufficient to calculate CIH from the scaled version of (1). For  
646 this reason, precipitation minus evaporation ( $P - E$ ) is used  
647 as a proxy for net convective heating where we interpret  
648  $P - E = 0$  as equivalent to a threshold of positive CIH.  
649 Figure 10a shows a multi-model average of the SST versus  
650  $P - E$  for the LGM and the P-IP depicting a 2 K change in  
651 the SST threshold for  $P - E = 0$ . Figure 10b presents a  
652 scatter plot of the cumulative positive  $P - E$  (i.e. the sum of  
653 the residual  $P - E$  where the residual is positive) in the  
654  $30^{\circ}\text{S}$ – $30^{\circ}\text{N}$  band for the LGM and P-IP periods. These rep-  
655 resent a proxy for the relative change in convective activity  
656 between the two periods. It is important to note that while the  
657 four models vary widely in their simulations of total positive  
658  $P - E$  (ranging from 1,800 to 8,000 mm in the P-IP) all  
659 the four models show larger magnitudes of total positive

660  $P - E$  (more precipitation exceeds evaporation) during the  
661 warmer P-IP compared to the colder LGM. The increase is  
662 about 6%. Figure 10c shows the change in area of positive  
663  $P - E$ . Three out of the four models show a slight reduction  
664 in area ( $<0.5\%$  in magnitude) in the P-IP compared to the  
665 LGM while one model shows a 2% increase. Ostensibly, the  
666 area is constant between the two climate epochs. Figure 10d  
667 summarizes the results of PMIP I analysis and indicates that  
668 the change in convective activity is about 5–10 times larger  
669 than the change in the size of the convective region.

670 Results using the PMIP II simulations are summarized in  
671 Fig. 11. The analysis includes six different models.<sup>2</sup> Fig-  
672 ure 11a shows the standardized average CIH in each  $0.5^{\circ}\text{C}$   
673 bin for the LGM, M-HP and the P-IP plotted against SST,  
674 relative to the P-IP. In other words, for each of the models,  
675 the SST versus CIH curves are plotted after subtracting  $T_H$  in  
676 the P-IP in the corresponding model from the SST values.

<sup>2</sup> NCAR CCSM, CNRM CM33, Hadley Center CM3M2, LASG 2FL01  
FGOALS, and MRI CGCM2.3.4. See <http://pmip2.lscce.ipsl.fr/> for 2FL02  
details. 2FL03



**Fig. 7** **a** Pentad evolution of the total tropical CIH binned every 0.5°C SST from 1950–1954 to 2000–2004 using the NCEP-NCAR reanalysis and the NOAA Extended Reconstructed SST v2. Different colours represent different pentads: cold colours represent pentads closely to 1950–1954 (shown in purple) with warm colours representing pentads closer to 2000–2004 (shown in red). The dotted box

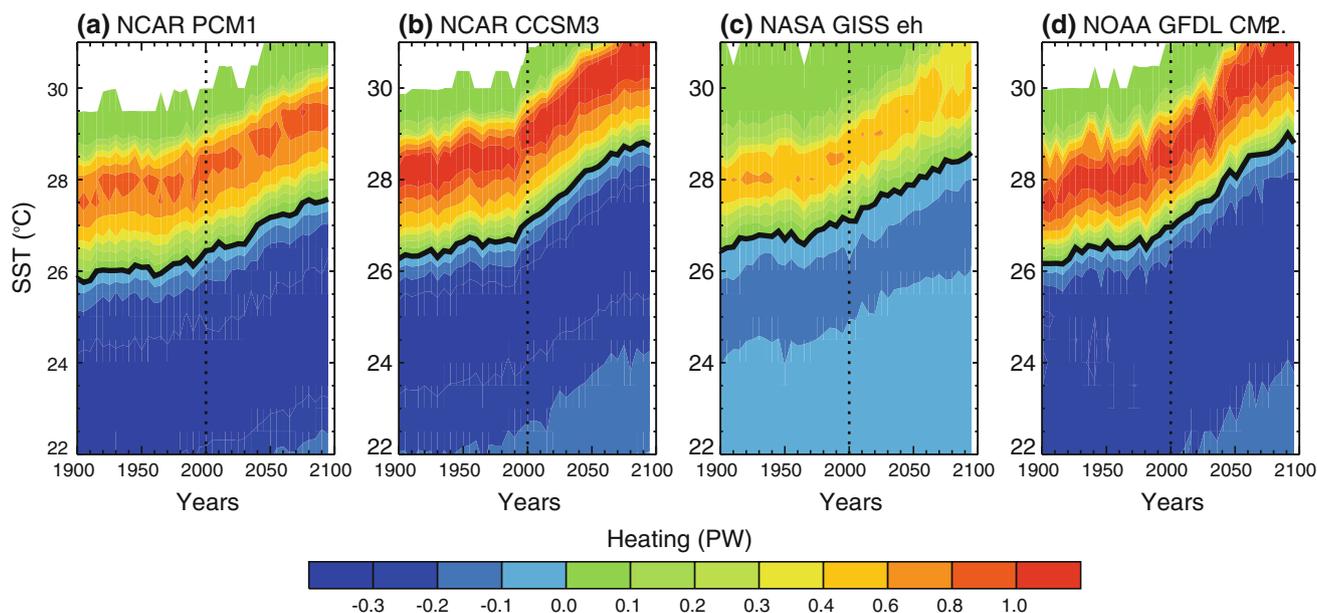
highlights the changes of the SST-threshold temperature  $T_H$  for atmospheric heating for different pentads. **d** Pentad evolution of  $T_H$  (black line), and net heating in the tropics (red line) from 1950–1954 to 2000–2004. **c** and **d** similar to **a** and **b**, respectively, but using ERA-40 reanalysis from 1960–1964 to 1995–1999

677 Note that all curves corresponding to the P-IP cross from  
 678 negative to positive CIH at zero relative SST. Also, the  
 679 standardization is performed by dividing the average CIH  
 680 values by the maximum absolute average CIH value in each  
 681 model run, thus obtaining values ranging from  $-1$  to  $1$ . The  
 682 average CIH is presented in a standardized manner and the  
 683 SST is shown relative to the corresponding  $T_H$  in the P-IP to  
 684 allow a direct comparison among the different climate  
 685 periods, leaving aside differences among model climatologies  
 686 due to their varying resolutions, differences in parameterizations  
 687 and biases. The variation in average temperature for a given  
 688 climate period among models could be as much as 1 K. PMIP II  
 689 simulations show an average 2 K difference in  $T_H$  between the  
 690 P-IP and the LGM, while during the M-HP the  $T_H$  is almost  
 691 the same as in the P-IP. The thresholds calculated from the  
 692 models are 24, 25.8 and 26°C for LGM, M-HP and P-IP  
 693 conditions, respectively, noting that the current threshold  
 694 temperature for convection is about 27°C (Fig. 3b).  
 695

Figure 11b shows the area of SST  $>26, 27$  and  $28^\circ\text{C}$  for the three different periods. The diagram also shows the SST thresholds  $T_H$  (right-hand scale). The results are similar to the PMIP I, CMIP3 and twentieth century observations. As temperature increases relatively from one period to another, the OWP also increases, but so does the convective threshold rendering the DWP almost constant in area. Also, both the PMIP I and PMIP II simulations suggest that the net CIH within the DWP is enhanced as temperature increases, again in agreement with twentieth and twenty-first century results.

**4 Constancy of the area of the dynamic warm pool (DWP)**

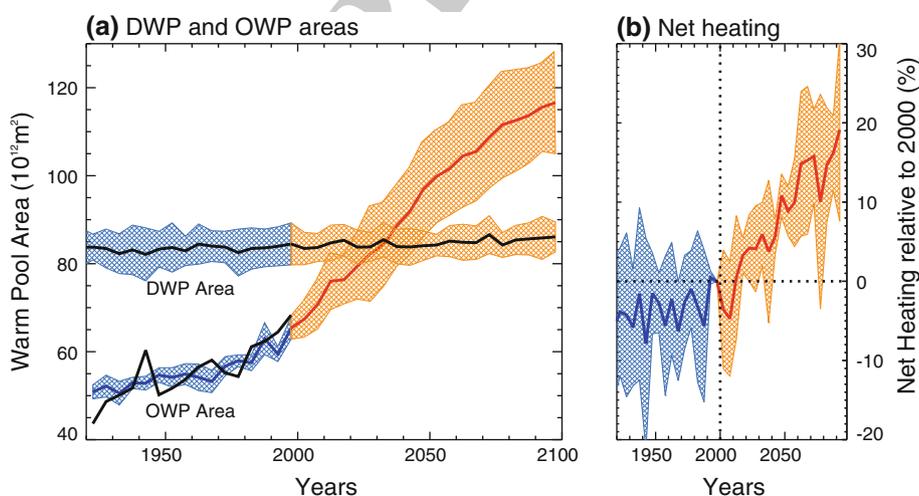
If the area of positive CIH in the tropical oceans remains constant as SST changes, the increased heating within the dynamic warm pool must be balanced by an increased



**Fig. 8** Evolution of CIH between 20°S and 20°N for the twentieth and the twenty-first centuries (A1B scenario) as a function of time and SST for the NOAA GFDL-CM2.1., NCAR PCM1, NCAR CCSM3

and NASA GISS-eh models. The atmospheric heating was integrated in 0.5°C SST bins. The *black line* corresponds to the threshold temperature,  $T_H$

**Fig. 9 a** Average evolution of the DWP area in the twentieth century and IPCC A1B (twenty-first century) scenarios for four different WCRP CMIP3 models (NCAR CCSM3, NCAR PCM1, NOAA GFDL 2.1 and NASA GISS eh). The DWP area does not include bias corrections. OWP area is also shown for the same scenarios for direct comparison. **b** Net heating over the tropical ocean for the twentieth century and the A1B CMIP3 scenarios relative to the 1995–1999 values



712 cooling in regions where  $SST < T_H$ . Similarly, in a cooler  
 713 environment (e.g., during the LGM) where the region of  
 714 CIH  $> 0$  possesses the same area, a similar compensation  
 715 must occur. We investigate this hypothesis using two  
 716 simple model prototypes.

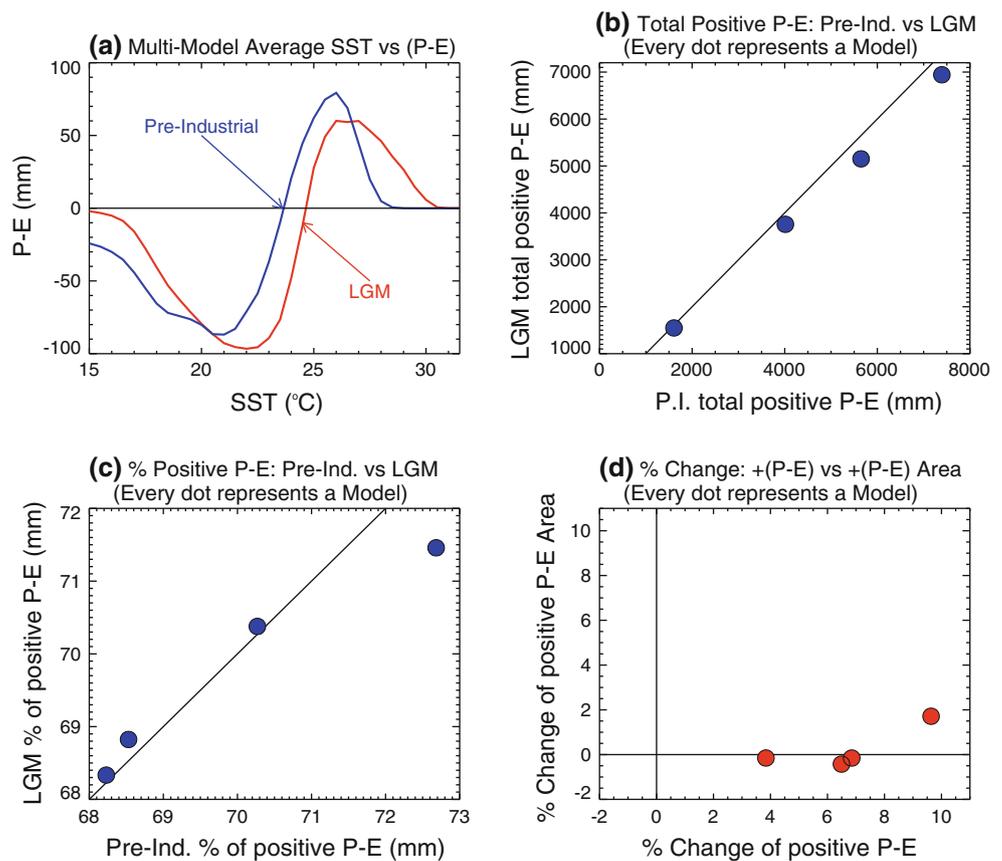
717 The hypothesis posed above suggests that the “vigour”  
 718 of the tropical circulation will increase as the average SST  
 719 of the tropics increase. But a number of studies have  
 720 reached an opposite conclusion: that as SST increases the  
 721 tropical circulations, especially the Walker Cell, reduces  
 722 intensity (e.g., Vecchi et al. 2006; Vecchi and Soden 2007).  
 723 We will return to this issue later noting that a number of  
 724 other papers (e.g., Zhou et al. 2011; Meng et al. 2011)

support a more vigorous tropical circulation, in accord with  
 the conclusions reached in the present study.

Consider a simple “two-box” model of the climate  
 system (Fig. 12a). The first box, corresponding to the area  
 of CIH  $> 0$ , or the DWP region, is dominated by rising  
 motion and convective latent heat release. The second box  
 encompasses the surrounding colder region, where air is  
 descending and cooling by long-wave radiative (LWR) lost  
 to space. In the cooling box the vertical convergence of  
 heat is given by the difference between the LWR entering  
 the column at the surface and the LWR exiting to space at  
 the top of the atmosphere (TOA). The LWR is given by the  
 Stephan–Boltzman law:

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**Fig. 10** Summary of analysis of the PMIP I models (NCAR CCM3, Canadian Climate Center model, GFDL, and UGAMP) for the LGM and P-IP. **a** Multimodel average of the SST versus  $P - E$  for LGM (blue) and P-IP (red). **b** Scatter plot of the cumulative positive  $P - E$  in the 30°S–30°N band for the LGM and P-IP. **c** Scatter plot of area of positive  $P - E$  for the LGM and P-IP. **d** Relative change in positive  $P - E$  and positive  $P - E$  area between the LGM and the P-IP



$$F_{net} = \varepsilon \sigma T^4 \quad (2)$$

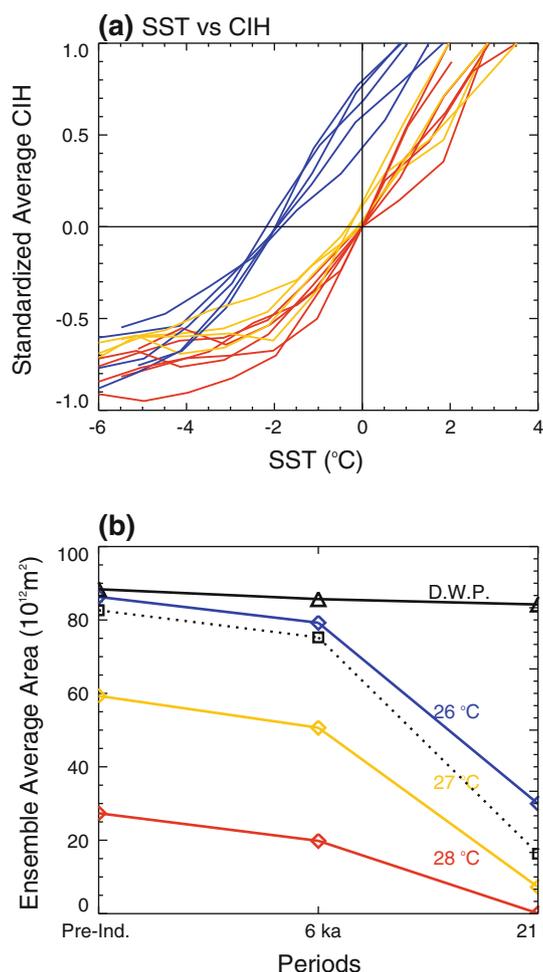
739 where  $\varepsilon$  is an emissivity and  $\sigma$  the Stefan–Boltzman  
 740 constant. In the warming box the vertical convergence of  
 741 heat is given by the difference of LWR at the surface and at  
 742 TOA plus the atmospheric heating within the column from  
 743 condensation ( $L_E$ ) where:

$$L_E = \rho_a L C_{DE} U q_s^* [(1 - R_H) + R_H L (T - T_a) / R_v T^2] \quad (3)$$

745 after Hartmann (1994). It is assumed, to first approximation,  
 746 that all evaporation from the warming and cooling boxes  
 747 condenses within the convective box. In (3),  $\rho_a$ ,  $L$ ,  $C_{DE}$ ,  $U$ ,  $q_s^*$ ,  
 748  $R_H$ ,  $R_v$ , and  $T_a$ , represent the density of air, the latent heat of  
 749 vaporization, the aerodynamic transfer coefficient for vapor,  
 750 the surface wind speed ( $\sim 5\text{--}10 \text{ ms}^{-1}$ ), the saturation specific  
 751 humidity at the surface, the relative humidity ( $\sim 70\%$ ), the gas  
 752 constant for water vapor and the air temperature.

753 Assume now that the two boxes are in equilibrium, with  
 754 convection and heating within the warm box balanced by  
 755 subsidence and cooling in the colder box. Note that surface  
 756 evaporation is a nonlinear function that increases expo-  
 757 nentially with SST through the Clausius-Clapeyron rela-  
 758 tionship. The radiative loss is also nonlinear, proportional  
 759 to the fourth power of temperature through the Stefan-  
 760 Boltzmann Law. Over a range of SST between 27 and  
 761 32°C, enclosing the LGM, the M-HP, the P-IP and the

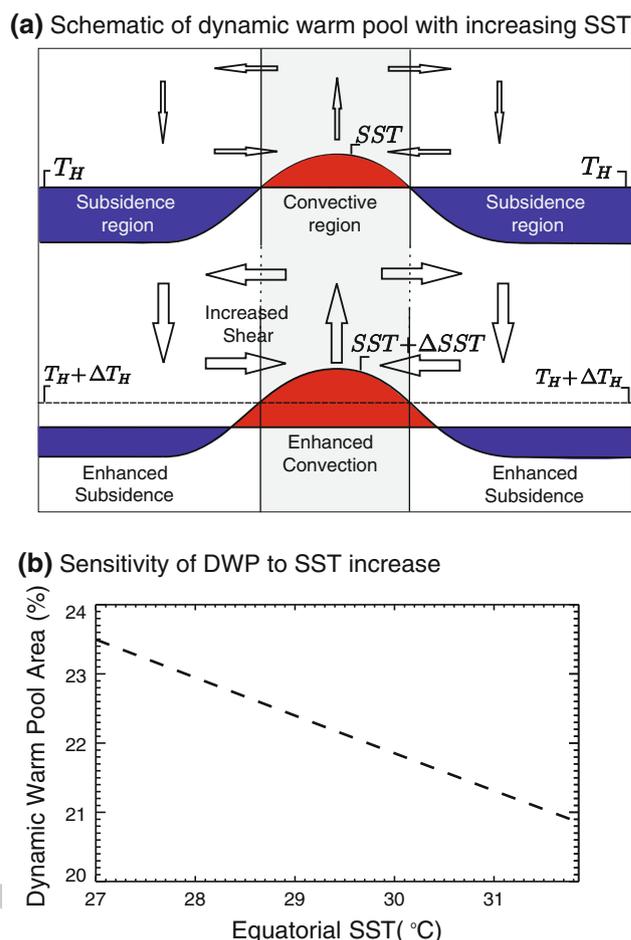
present climate as well as what may be expected reason- 762  
 ably in the future, the area of the dynamic warm pool 763  
 ranges from 23.5 to 21% of the surface of the idealized 764  
 planet (Fig. 12b) for a homogeneous SST increase in both 765  
 the heating and cooling boxes. These values are similar to 766  
 the  $\sim 25\%$  coverage of convection over ocean found in the 767  
 CMIP3 models and suggested much earlier by the analysis 768  
 of Gray (1973) using satellite observations. Thus, for a 769  
 homogeneous SST increase the area of the DWP will 770  
 essentially remain constant, with a tendency to decrease 771  
 slightly as SST increases. This near-constancy suggests 772  
 that the sensitivity of radiative cooling to space to changes 773  
 in SST (i.e.,  $dF_{net}/d(\text{SST})$  from Eq. 2) must be almost equal 774  
 (and opposite) to the sensitivity of latent heat release in 775  
 within a column to SST (i.e.,  $dL_E/d(\text{SST})$  from Eq. 3). To 776  
 explore this finding further, experiments were conducted 777  
 using a relatively simple two-level, nonlinear zonally 778  
 symmetric model with a full hydrological cycle (Webster 779  
 and Chou 1980a, b; Webster 1983). The model is adapted 780  
 to represent an aqua-planet (i.e. no land). The atmospheric 781  
 component is a primitive equation model extending from 782  
 pole to pole, with explicit calculation of latent heating due 783  
 to regional convective effects and large-scale convection 784  
 depending on moisture availability and equivalent potential 785  
 temperature following Ooyama (1969) and Anthes (1977). 786  
 Surface fluxes are calculated at the surface relative to the 787



**Fig. 11** Summary of analyses of the PMIP II models (NCAR CCSM, CNRM CM33, Hadley Center CM3M2, LASG FGOALS, and MRI CGCM2.3.4) for the LGM, M-HP and P-IP. **a** SST versus average CIH in 0.5°C bins for the LGM (blue), M-HP (orange) and P-IP (red). The average CIH is presented in a standardized manner and the SST is shown relative to the corresponding  $T_H$  in the P-IP. **b** Area of SST > 26, 27 and 28°C as well as the DWP area for the three different periods (black line). The diagram also shows the SST thresholds  $T_H$  (dotted line)

788 state of the evolving model atmosphere and the underlying  
 789 SST. The radiation scheme is configured as a two-stream  
 790 model following Stephens and Webster (1979, 1984).

791 The atmospheric component of the model is run in  
 792 uncoupled mode to equilibrium for different SST sinusoidal  
 793 configurations. Each SST forcing has the same latitudinal  
 794 distribution so that there is no change in meridional  
 795 gradient but with different magnitudes due to a homoge-  
 796 nous increase of surface temperature. Figure 13a shows  
 797 the different SST equatorially symmetric configurations  
 798 plotted in steps of 1°C from 26 to 38°C. Figure 13b–d  
 799 show the resulting annual average vertical velocity, vertical  
 800 wind shear in the meridional direction and 250 hPa tem-  
 801 perature. For easier interpretation, the fields are coded in



**Fig. 12 a** Schematic diagram of the “two-box” model containing a convective and a subsidence region. The *top panel* represents the current state of the climate where it is assumed that the convective ( $SST > T_H$ ) and subsidence regions are in dynamical and thermodynamical balance. The *bottom panel* represents the two-box model after a homogeneous increase in SST in both boxes. Here, the heating threshold SST increases ( $T_H + \Delta T_H$ ), but the convective region remains constant. The balance and constancy of in the climate state is achieved by both enhanced convection and subsidence. **b** Dynamic warm pool (convective region) size relative to the surface area of the planet in the simple two-box model

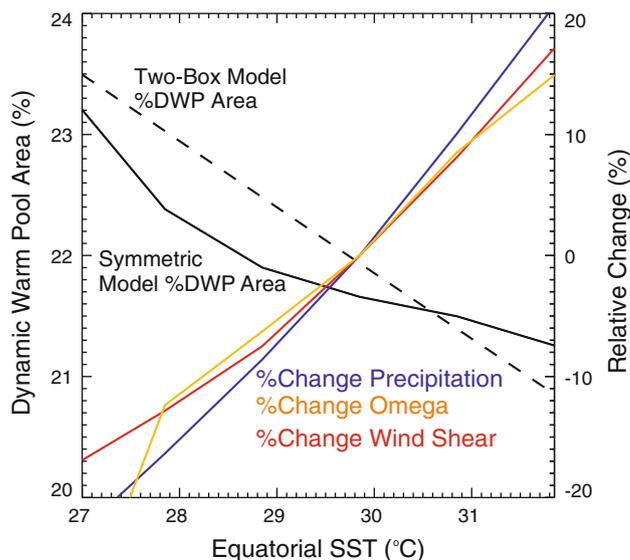
802 the same colors as the SST configuration shown in  
 803 Fig. 13a. The model simulations result in a zone of narrow  
 804 convection at the equator and subsidence in both hemi-  
 805 spheres poleward of about 15°N and S. Both equatorial  
 806 convection and off-equatorial subsidence increase as the  
 807 SST forcing increases in magnitude, enhancing the  
 808 meridional circulation (Fig. 13b) and precipitation. That is,  
 809 as the base SST increases, the tropical circulations become  
 810 more vigorous. Also, in accordance with the increased  
 811 vigour, the meridional vertical wind shear also increases at  
 812 the edges of the convective cells (Fig. 13c). The upper  
 813 tropospheric temperature (250 hPa) also increases. The  
 814 increase of the 250 hPa temperatures is greater at low

815 latitudes because of the non-linearity of moist processes  
 816 (Fig. 13d). Note that while the range of SST variations at  
 817 the equator is  $\sim 10^\circ$ , the range at the 250 hPa is  $\sim 20^\circ$ .  
 818 Thus, the dry static energy of the model increases with SST  
 819 but the moist static energy decreases, or becomes more  
 820 unstable.

821 Figure 14 summarizes the changes in the relative area of  
 822 the convective versus the non-convective regions of the  
 823 tropics for the box model and the aqua-planet model. The  
 824 size of the convective area in the zonally symmetric aqua-  
 825 planet behaves remarkably similar to the two-box model,  
 826 staying relatively constant, and also with a tendency to  
 827 slightly decrease with increasing equatorial SST. In addition,  
 828 results indicate that the total precipitation, the magni-  
 829 tude of the average vertical velocity in the convective  
 830 region and the vertical wind shear at the edge of the  
 831 dynamic warm pool increase considerably with SST  
 832 (Fig. 14).

833 **5 Summary and conclusion**

834 In the tropics, the distribution and magnitude of the domi-  
 835 nant latent heat release associated with deep convective  
 836 activity is, at any given time, strongly associated with the  
 837 distribution of SST. In the 1979–2001 period the net  
 838 atmospheric heating in the  $30^\circ\text{S}$ – $30^\circ\text{N}$  belt occurs where  
 839  $\text{SST} > 27^\circ\text{C}$ , and the spatially averaged heating increases  
 840 nonlinearly with SST (Fig. 3b). We argue subsequently  
 841 that this region of net positive atmospheric heating could  
 842 be referred to as the dynamic warm pool (DWP) as it is  
 843 determined not only by the spatial distribution of the SST

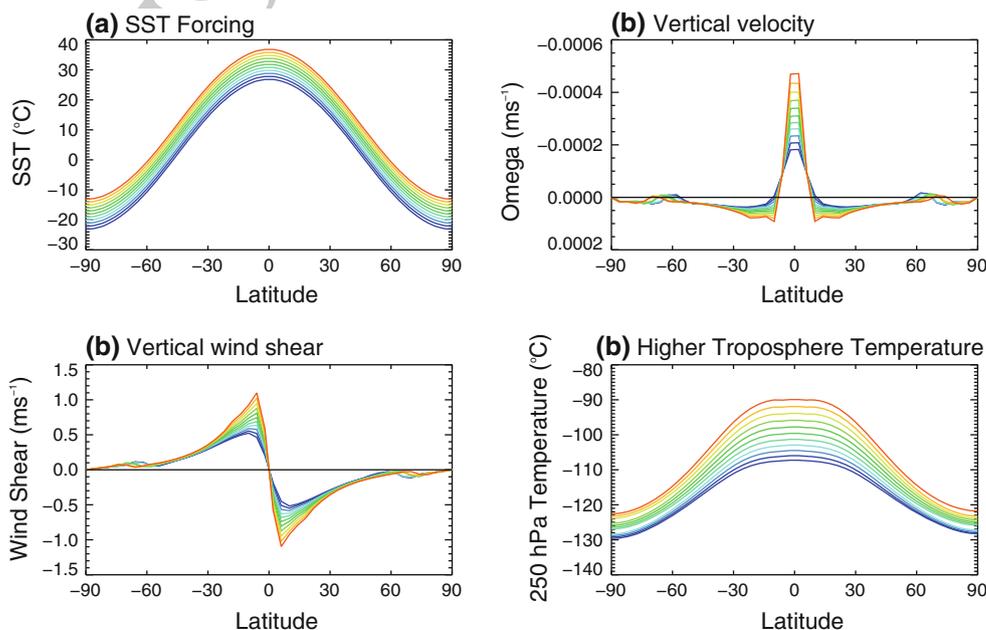


**Fig. 14** Size of the DWP relative to the surface area of the planet in a simple two-box model (dashed black line, same as Fig. 12b for comparison), and in the zonally symmetric aqua-planet (continuous black line). Changes of total precipitation (blue line), magnitude of average vertical velocity within the convective region (red line), and wind shear at the edge of the convective region (orange line) are plotted relative to the present (SST  $\sim 29^\circ\text{C}$ )

844 but also the large scale atmospheric circulation. In other  
 845 words, rather than a static definition set by a constant  
 846 temperature, the climatically active warm pool is defined  
 847 dynamically by the large-scale moist coupled ocean–  
 848 atmosphere system.

849 The nature of the relationship between local and remote  
 850 SST with local convection during the past, present and

**Fig. 13** **a** Configurations of SST forcing for the zonally symmetric nonlinear global primitive equation model varying the equatorial temperature in steps of  $1^\circ$  from 26 to  $38^\circ\text{C}$ , **b** simulated vertical velocity, **c** simulated meridional vertical wind shear, and **d** 250 hPa temperature for each of the SST configurations. The matching colours in all four panels correspond to the same experiment



851 future climate states are determined by examining the  
 852 evolution of an OWP arbitrarily defined by SST >28°C and  
 853 the DWP, as well as the magnitude of the net heating  
 854 within the tropical DWP. Observations and the PCM,  
 855 CMIP3, and PMIP model simulations are used to study the  
 856 long-term variability of both forms of warm pool during  
 857 the Last Glacial Maximum (21 ky BP) and the M-HP  
 858 period (6 ky BP) compared to P-IP conditions, and their  
 859 state during the current era and the evolution that may  
 860 occur during the next century. These different climate  
 861 epochs correspond to vastly different climate states and  
 862 differing forcing factors.

863 During the twentieth century tropical SST has increased  
 864 by about 0.8°C in all tropical oceans accompanied by a  
 865 steady 70% expansion of the OWP area, spread uniformly  
 866 throughout the annual cycle. CMIP3 models are used to  
 867 assess the twenty-first century changes in OWP and DWP.  
 868 The models are first bias-corrected using twentieth century  
 869 simulations. In general, all models simulate the observed  
 870 twentieth OWP expansion remarkably well although  
 871 slightly underestimating the steep increase after the 1970s.  
 872 Multi-model bias-corrected twenty-first century projections  
 873 indicate that the OWP and the area of SST >26.5°C  
 874 increase for the COMMIT scenario (i.e., CO<sub>2</sub> concentra-  
 875 tions held at 2,000 levels) about 10–20% compared to the  
 876 2000–2004 value, while for scenarios A1B and A2 the  
 877 OWP area increases 70 and 90%.

878 Analysis of the evolution of CIH in the tropics during  
 879 the last 50 years reveals a non-stationary relationship with  
 880 SST, with CIH profiles shifting to higher SSTs as the SST  
 881 increases (Fig. 8). CMIP3 simulations behave similarly to  
 882 observations during the twentieth century and project a  
 883 further shift of the CIH profiles to higher SSTs during the  
 884 twenty-first century. The shift in CIH implies an increase in  
 885 SST convective threshold. In other words, the transition  
 886 between different convective regimes (e.g., the onset of  
 887 deep convection and increases in area averaged CIH) while  
 888 dependant on local SST in the short and medium terms, is  
 889 not stationary relative to SST. As a result, the DWP area  
 890 does not vary significantly in the long-term, covering about  
 891 25% of the global ocean. That is, while the SST and the  
 892 OWP area increase sharply with time, the area of atmo-  
 893 spheric convective heating in the tropics remains remark-  
 894 ably constant over the range of climate observed in the  
 895 twentieth century and the simulated twenty-first century as  
 896 well as the M-HP and the LGM. In essence, with the  
 897 warming of SST during the twentieth century, the obser-  
 898 vational and model results point to an intensification of the  
 899 hydrological cycle as suggested by Trenberth (1999) with  
 900 intensified latent heat release within the DWP. These  
 901 concepts also allow an interpretation of the Lindzen and  
 902 Nigam (1987) results but for a changing climate. Lindzen  
 903 and Nigam pointed out the importance of the gradient of

SST in driving the tropical circulation. If the gradient of  
 SST were to remain the same as SST universally increases,  
 one would expect an increase in mass convergence from  
 Clausius-Clapeyron considerations consistent with the  
 results of the simple models presented in Sect. 4.

Paleoclimate simulations (PMIP I and II) for the LGM  
 and the M-HP are in agreement with the results from  
 CMIP3 simulations and twentieth century observations,  
 with enlargement of the OWP as temperature increases  
 relatively from one period to another, and a shift in the  
 convective SST threshold rendering the DWP almost  
 constant in area. The OWP during the M-HP and the P-IP  
 appears considerably larger than during the LGM. But, the  
 LGM convective threshold is about 2°C lower than either  
 the M-HP or the P-IP while maintaining a nearly constant  
 DWP area (Fig. 11). With a lower threshold temperature,  
 the integrated heating in the LGM DWP is considerably  
 smaller than the other periods considered (Fig. 11a) lead-  
 ing to a hydrological cycle of probably reduced vigour.  
 This reduction is consistent with the relative aridity  
 occurring in the tropical regions during the LGM (Webster  
 and Stretten 1978; Pinot et al. 1999).

The magnitude of CIH within the DWP increases sub-  
 stantially in the different twenty-first century model sim-  
 ulations accompanying the projected SST rise with an  
 estimated increase of about 20% in heating relative to  
 present day values. This increase implies an enhancement  
 of the global hydrological cycle, with an intensified cir-  
 culation, both in the ascending and descending branches,  
 resulting in a more vigorous vertical shear. In fact, analysis  
 of two simple climatic and atmospheric circulation models  
 suggest that, in a range of temperatures enclosing the  
 present climate and what may be expected reasonably in  
 the future, the increase heating within the warm pool is in  
 near balance with increase cooling outside following an  
 increase in global SST and leading to a nearly constant  
 dynamic warm pool area.

One of the major conclusions from this study, that the  
 tropical circulation has become more vigorous during the  
 period of warming since the P-IP, with expectations of  
 further enhancement in the twenty-first century is contrary  
 to a number of studies. For example, Vecchi et al. (2006),  
 Held and Soden (2006) and Vecchi and Soden (2007),  
 using a reanalysis data set, as well as suites of AMIP3  
 models, argue that the Walker and Hadley circulations  
 have reduced their intensity and will continue to reduce  
 their intensity in a world of warming SST. Held and Soden  
 (2006) have developed a heuristic thermodynamic argu-  
 ment to explain increases in lower tropospheric moist  
 (accompanying SST increases), increased precipitation but  
 lower mass flux. Mitras and Clement (2005, 2006) note that  
 reanalysis products have indicated accelerations of both the  
 Hadley and Walker Circulations during the last

30–40 years differing substantially from the Vecchi et al. (2006) conclusions. They also note that reanalysis products indicate a cooling of the mid-upper troposphere during the late twentieth century, contrary to the models that suggest generally a warming in the mid-upper troposphere. These differences in warming may explain differences in static stability that may help explain a reduction or acceleration of the large-scale tropical circulations. Zhou et al. (2011) use data sets independent of the reanalysis products (the Global Precipitation Climatology Program: GPCP, Adler et al. 2003, and the International Satellite Cloud Climatology Project: ISCCP, Rossow and Schiffer 1999) to determine the changes occurring in the Walker and Hadley cells. In agreement with Mitas CM and Clement (2005, 2006) they find an intensification of the major tropical circulations. In addition, Zhou et al. (2011) find increases in outgoing longwave radiation in the subsident regions of the tropics in accord with a warming SST and increased moisture in the boundary layer. These results are consistent with the results of our simple models presented in Sect. 4.

Meng et al. (2011) note the importance of determining SST gradients accurately both in observations and simulations. They suggest that warming of tropical SSTs has taken place asymmetrically and that if not carefully observed or modelled, errors could be made in determining trends in circulations. They note, especially inter-basin SST gradients between the poorly observed Indian Ocean and the western Pacific Ocean. Consistent with the recent papers Shin and Sardeshmukh (2011) and Shin et al. (2010), Meng et al. (2011) note the importance of models simulating the regional tropical SST patterns with fidelity in order to ascertain both present climate or future changes.

Potentially, an intensified hydrological cycle could have important thermodynamical and dynamical implications in tropical and global climate, ranging from the location and amount of convection and rainfall during past and future climates to the vigor of teleconnection patterns between the tropics and higher latitudes, and up-scale and down-scale transfer of energy. For example, the near constancy in DWP area, together with the intensified CIH and vertical shear are useful in exploring why observations and model simulations suggest a shift in the hurricane distributions towards more intense, yet fewer storms (Bender et al. 2010; Knutson et al. 2010). This particular issue will be addressed in a companion paper. It is argued that the increasing magnitude of the atmospheric heating within the dynamic warm pool is in accord with an increase in the average intensity of tropical cyclones, as disturbances have higher likelihood of intensification in large-scale moist-convective high-CIH environments (Agudelo et al. 2010). However, accompanying the increased heating within the dynamic warm pool is an increase of mass flux in and out of the region of net heating, increasing the vertical wind shear at

the peripheries of the dynamics warm pool (Fig. 14). This increase in vertical wind shear may provide a constraint on cyclone formation, potentially limiting the number of tropical cyclones, even as their average intensity increases.

**Acknowledgments** The research was funded by the National Science Foundation Climate Division Grant ATM-0531771. The authors would like to acknowledge colleagues and members of our research group (Paula Agudelo, Judith Curry, Hye-Mi Kim, Violeta Toma, Hai-Ru Chang, James Belanger) for many discussions regarding the results of the study. We appreciate the advice of Sumant Nigam for discussions and advice about vertically integrated heating. We acknowledge Program for Climate Model Diagnosis and Intercomparison (PCMDI) and the WCRP's Working Group on Coupled Modelling (WGCM) for making available the WCRP CMIP3 multi-model dataset. Support of this dataset is provided by the Office of Science, U.S. Department of Energy (DOE). Data from the Parallel Climate Model runs with different anthropogenic and natural forcing were obtained from the U.S. DOE Couple Climate Model Data Archives ([http://www.nerse.gov/projects/gcm\\_data/](http://www.nerse.gov/projects/gcm_data/)) through the Earth System Grid (ESG, <http://www.earthsystemgrid.org/>). We also acknowledge ECMWF for making the 1979–2001 CIH available to us for this study. ERA-40 reanalysis data used in this study was also supplied by ECMWF. NCEP Reanalysis Derived data provided by the NOAA/OAR/ESRL PSD (Boulder, Colorado, USA) from their Web site at <http://www.cdc.noaa.gov/>.

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